

Effective Planetary Defense using Directed Energy

DE-STARLITE

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Abstract— We show that Directed Energy (DE) systems offer the potential for true planetary defense from small to km class threats. Directed energy has evolved dramatically recently and is on an extremely rapid ascent technologically. It is now feasible to consider DE systems for threats from asteroids and comets. DE-STAR (Directed Energy System for Targeting of Asteroids and exploRation) is a phased-array laser directed energy system intended for illumination, deflection and compositional analysis of asteroids [1]. It can be configured either as a stand-on or a distant stand-off system. A system of appropriate size would be capable of projecting a laser spot onto the surface of a distant asteroid with sufficient flux to heat a spot on the surface to approximately 3,000 K, adequate to vaporize solid rock. Mass ejection due to vaporization creates considerable reactionary thrust to divert the asteroid from its orbit. DE-STARLITE is a smaller stand-on system that utilizes the same technology as the larger standoff system, but with a much smaller laser for a dedicated mission to a specific asteroid. DE-STARLITE offers a very power and mass efficient approach to planetary defense. As an example, a DE-STARLITE system that fits within the mass and size constraints of the Asteroid Redirect Mission (ARM) system in a small portion of the SLS block 1 launch capability is capable of deflecting an Apophis class (325 m diameter) asteroid with sufficient warning. A DE-STARLITE using the full SLS block 1 launch mass can deflect any known threat. We propose a logical approach to planetary defense is to pre-deploy DE-STARLITE systems in LEO or GEO rather than start a build when a threat arises. In the times when the system is not being used for planetary defense it can be used for many other tasks including orbital debris removal. Pre-deployment allows for rapid repose to threats and is a far superior approach to waiting for a threat before a mission start. We compare DE-STARLITE to other deflections possibilities including impactor and ion beam deflection missions with the same launch mass. DE approaches are far superior to their deflection capability.

Keywords—DE-STAR; DE-STARLITE; Planetary Defense; Directed Energy; Laser Phased Array

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1. INTRODUCTION

This paper first introduces the motivation behind implementing a directed energy planetary defense system as it acknowledges the need for planetary defense and explains the benefit of utilizing laser ablation of an asteroid over any alternative method to impart a deflecting force on the threat. The general proposed system is called DE-STAR, for Directed Energy System for Targeting of Asteroids and exploRation. The specific mission, detailed in Section 2 of this paper, is called DE-STARLITE—a dedicated stand-on mission that utilizes much of the same technology but is fundable and feasible on a shorter time scale due to its smaller scope. Orbital deflection models have been developed to understand the orbital deflection capabilities of such a system, as is detailed in Section 3 of this paper.

Asteroid Impact Threat

Asteroid impacts pose a continual threat to modern

civilization. On 15 February 2013, an asteroid penetrated the atmosphere over Chelyabinsk, Russia entering at an angle of approximately 18°, and releasing energy equivalent to 570 ± 150 kt TNT [2]. For comparison, the nuclear weapon that was detonated approximately 509 m above the ground in Hiroshima, Japan yielded approximately 12.5 kt TNT [3]. The main airburst over Chelyabinsk occurred at an approximate altitude of 30 km and created a shock wave strong enough to shatter windows out to a distance of 120 km from the meteorite's track, injuring over 1,200 people in Chelyabinsk city and hundreds more in nearby towns and rural areas [2]. Had the asteroid approached from a higher angle, more serious damage would be anticipated from higher concentration of the impact energy on the ground.

Sixteen hours after the meteorite struck near Chelyabinsk, the 45 m diameter asteroid 2012 DA14 approached to within 27,743 km of Earth's surface—inside the orbit of geosynchronous satellites. If DA14 were to strike Earth, it would deliver approximately 7.2 Mt TNT [4]. Although the Chelyabinsk meteorite and DA14 arrived at or near Earth on the same day, the two objects were not linked to each other, coming from completely unrelated orbits. That two such seemingly improbable events could occur within hours of each other serves as a stark reminder that humanity is continually at risk of asteroid impact.

Asteroids at least the size of DA14 (~50 m diam.) are expected to strike Earth approximately every 650 years, while objects at least the size of the Chelyabinsk impactor (~20 m diam.) are expected to strike Earth approximately every 100 years [4]. Larger objects also pose a severe threat, as the total kinetic energy associated with an impact of a 100 m asteroid is equivalent to approximately 85 Mt TNT, and that of the well-known 325 m threat, Apophis, is approximately 3.2 Gt TNT [4]. Thus, effective mitigation strategies are imperative to ensure humanity's continuity and future advancement.

Mitigation Methods

Several concepts for asteroid deflection have been described, which can be broadly generalized into six distinct strategies.

- (1) Kinetic impactors, with or without explosive charges: An expendable spacecraft is sent to intercept the threatening object. Direct impact would modify the object's orbit through momentum transfer. Enhanced momentum transfer can be accomplished using an explosive charge, such as a nuclear weapon [5], [6], [7], [8].
- (2) Gradual orbit deflection by surface albedo alteration: The albedo of an object could be changed using paint [9], mirrors [10].
- (3) , sails [11], *etc.* As the albedo is altered, a change in the object's Yarkovsky thermal drag would gradually shift the object's orbit.
- (4) Direct motive force, such as by mounting a thruster directly to the object: Thrusters could include chemical propellants, solar or nuclear powered electric drives, or ion engines [12]. Such methods, including ion beam deflection (IBD), require much greater mission mass than does the

laser ablation method, as proposed for the DE-STARLITE mission [13].

- (5) Indirect orbit alteration, such as gravity tractors: A spacecraft with sufficient mass would be positioned near the object, and maintain a fixed station with respect to the object using onboard propulsion. Gravitational attraction would tug the object toward the spacecraft, and gradually modify the object's orbit [14], [15].

- (6) Expulsion of surface material, *e.g.* by robotic mining: A robot on the surface of an asteroid would repeatedly eject material from the asteroid. The reaction force from ejected material affects the object's trajectory [16].

- (7) Vaporization of surface material: Similar to robotic mining, vaporization on the surface of an object continually ejects the vaporized material, creating a reactionary force that pushes the object into a new path. Vaporization can be accomplished by solar concentrators [17] or by lasers [18] deployed on spacecraft stationed near the asteroid, the latter of which is proposed for the DE-STARLITE mission (Section 3). During laser ablation, the asteroid itself becomes the "propellant"; thus a very modest spacecraft can deflect an asteroid much larger than would be possible with a system of similar mission mass using alternative techniques.

2. DE-STAR

The DE-STAR concept is envisioned as an orbiting system consisting of a modular array of phase-locked lasers powered by photovoltaics [1]. The multi-purpose system is capable of planetary defense against asteroids that are projected to collide with the Earth. Laser ablation of the asteroid imparts a deflecting force on the target in order to mitigate the risk of impact. The laser produces a spot on the target that heats the surface at the spot to a temperature great enough to vaporize all known constituent materials—approximately 3,000 K. The vaporization consequently creates a reactionary force that diverts the asteroid. Recent advances in photonics make a scientific discussion of directed energy planetary defense feasible whereas even 10 years ago it was close to science fiction. High power lasers are capable of delivering sufficient energy density on a target to melt and vaporize any known material. Laser machining and welding are commonplace in industry, where even refractory metals are directly machined or joined with lasers. Scaling of laser technology has spurred development of directed energy systems that are capable of delivering high energy density on distant targets. Recent developments have resulted in conversion of electrical to photon efficiencies of close to 50% with powers in excess of 1 kW per (handheld) unit. Additionally, and critical for this program, such devices can be phased locked. This field is rapidly changing and even more efficient devices with higher power density will be available in the near future. This allows us to contemplate directed energy systems for large scale deployment. Inside the Earth's atmosphere, directed energy systems are hindered by atmospheric fluctuations of the coherent beam. A directed energy system deployed above the atmosphere could project a beam

through space unfettered by atmospheric interference and thus allows us to design systems that are essentially diffraction limited as the interplanetary medium (IPM) is extremely tenuous and does not affect the laser beam significantly. The system consists of a large array of phase-locked modest power laser amplifiers. By controlling the relative phases of individual laser elements, the combined beam can be directed to a distant target. Lasers are powered by solar photovoltaics of essentially the same area as the laser array. By increasing the array size we can both reduce the spot size due to diffraction and increase the power. This dual effect allows us to vaporizing elements on the surface of asteroids at distances that are significant compared to the solar system. By raising the flux (W/m^2) on the target asteroid to a sufficiently high level we can begin direct evaporation of the asteroid at the spot. This has two basic effects. Firstly, we directly begin to evaporate the asteroid and given sufficient time, a threatening asteroid could be totally vaporized before hitting the Earth. Secondly, evaporation at the spot causes a back reaction on the asteroid from the vaporization plume which acts as a rocket and thus the asteroid can be deflected. Since DE-STAR is a phased array consisting of a very large number of elements it can simultaneously be used for multiple purposes and is intrinsically a multi-tasking system. Fig. 1 depicts an orbiting DE-STAR system simultaneously engaged in both evaporating and deflecting a large asteroid as well as powering and propelling a spacecraft. The system consists of an array of phase-locked lasers. By controlling the relative phases of individual laser elements, the combined beam can be directed to a distant target. Lasers are powered by a solar panel of effectively the same area as the laser array. A DE-STAR of sufficient size would be capable of vaporizing elements on the surface of asteroids. Given sufficient time, a threatening asteroid could be vaporized, deflected or disintegrated prior to impacting Earth. The ability to direct energy onto a distant target renders DE-STAR capable of many functions. Asteroid interrogation may be possible by viewing absorption lines as the heated spot is viewed through the ejected vapor plume. Photon pressure can be used to accelerate (and decelerate) interplanetary spacecraft, among many other possibilities.



Figure 1. Left: Concept diagram of an orbiting DE-STAR engaged in multiple tasks including asteroid diversion, composition analysis and long range spacecraft power and propulsion. **Right:** Visualization with relevant physical phenomenon included at a flux of about $10 \text{ MW}/\text{m}^2$. Plume density is exaggerated to show ejecta. Asteroid diameter is about that of Apophis (325 m) relative to the laser beam diameter (30 m). Target is at 1 AU.

As this is a modular system we classify each DE-STAR by the log of its linear size, thus a DE-STAR 1 is 10 m, DE-STAR 2 is 100 m, *etc.* A DE-STAR 4 system will produce a reaction thrust comparable to the Shuttle SRB on the asteroid due to mass ejection and thus allow for orbital diversion of even larger asteroids, beyond several km in diameter, thus allowing for protection from every known asteroid threat. Smaller systems are also extremely useful. For example, a DE-STAR 2 (100 m array) would be capable of diverting volatile-laden objects 100 m in diameter by initiating engagement at $\sim 0.01\text{-}0.5 \text{ AU}$ ($\text{AU} = \text{Astronomical Unit} = \text{mean distance from Earth to Sun} \sim 1.5 \times 10^{11} \text{ m}$). Smaller objects could be diverted on shorter notice. The phased array configuration is capable of creating multiple beams, so a single DE-STAR of sufficient size could engage several threats simultaneously, such as a Shoemaker-Levy 9 scenario on Earth. An orbiting DE-STAR would also be capable of a wide variety of other functions. Narrow bandwidth and precision beam control would aid narrow search and ephemeris refinement of objects identified with wide-field surveys. Propulsion of kinetic or nuclear tipped asteroid interceptors or other interplanetary spacecraft is possible using the "photon rail gun" mode from direct photon pressure on a spacecraft, propelling (for example) a wafer scale spacecraft to $c/4$ in 10 minutes to reach the nearest stars in about 15 years, a 100 kg craft to 1 AU in 3 days or a 10,000 kg craft to 1 AU in 30 days. Vaporization and de-orbiting of debris in Earth orbit could be accomplished with a DE-STAR 1 or 2 system. DE-STAR 3 and 4 arrays may allow standoff interrogation of asteroid composition by observing absorption lines in the blackbody spectrum of a vaporizing surface spot. There are a number of other applications as well, including downlink power via mm, microwave or laser—the so called Space Power System mode. The system is a standoff planetary defense system that is always ready when needed and no dedicated mission is needed for each threat as is the case with other proposed mitigation methods.

3. DE-STARLITE MISSION

While the larger DE-STAR system remains a long term goal, DE-STARLITE is a more feasible and fundable mission as it is a smaller, stand-on version of the larger standoff system. DE-STARLITE is designed to be sent on a spacecraft with a 1 m to 4.5 m diameter array, to arrive nearby a Near-Earth Asteroid (NEA) and deflect it from its potentially hazardous trajectory. The laser array is essentially the same as for the DE-STAR program but vastly smaller. A secondary approach with a lower risk potential fallback is a close-packed focal plane array of fiber lasers [13]. DE-STARLITE is made possible with high-power solar electric propulsion (SEP). PV panels will be stowed for launch and will deploy upon reaching low-Earth orbit (LEO) to provide a required 100 kW electrical power from two 15 m diameter ATK MegaFlex panels. Even larger power is possible within the launch mass and shroud sizes available. The system will utilize ion engines (detailed

below) to propel the spacecraft from LEO to an NEA, as proposed in JPL’s ARM program [13]. The system aims to stay within the same mass and launch constraints as ARM and use much of the same propulsion technology. The laser efficiency determines the laser power obtained from the PV arrays; 35 kW of laser power would be produced at 35% efficiency, 50 kW at 50%, and 70 kW at 70%. The 35 kW estimate is based on the current efficiency (35%) of existing technology of the baseline Ytterbium laser amplifiers and thus provides for the worst case, while the 50 and 70 kW estimates are based on feasible technological improvement within the next 5-15 years. For example, 50% efficiency looks readily achievable within less than 5 years. A passive cooling radiator with z-folded arrays will be used to reject waste heat and maintain the temperature at near 300 K. Conceptual drawings of the system and payload are shown in Fig. 2 and Fig 3.

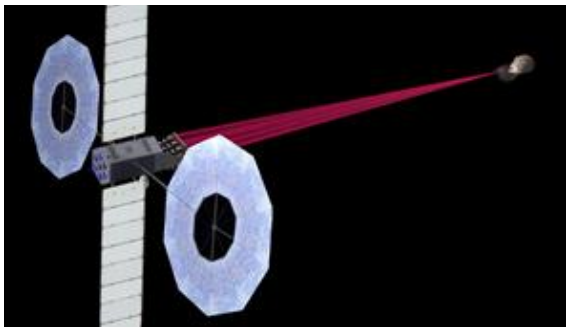


Figure 2. Artistic rendering of a deployed DE-STARLITE spacecraft deflecting an asteroid. The spacecraft is outfitted with two 15 m diameter MegaFlex PV Arrays, a z-folded radiator deployed up and down, a laser array mounted on a gimbal at the front, and ion engines at the back. From Kosmo *et al.* [13].

The PV panels are currently scalable to about 450 kW per pair and have a mass per unit power of about 7 kg/kW. The minimum flux on target requirement is set by the material properties. We have focused our work on the worst case of high temperature materials that require spot temperature of 2,000-3,000 K for efficient mass ejection. This is discussed in detail in a series of papers our group has published. An example of a 3D simulation for a typical rocky material is shown. Surface flux above 10 MW/m² is sufficient to efficiently ablate most materials of interest.

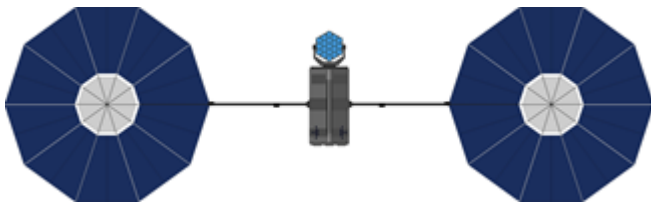


Figure 3. Conceptual design of the deployed spacecraft with two 15 m PV arrays that produce 50 kW each at the beginning of life for a total of 100 kW electrical, ion engines at the back, and the laser array pointed directly at the viewer. A 2 m diameter laser phased array is shown with 19 elements, each of which is 1-3 kW optical output. From Kosmo *et al.* [13].

Radiators

Thermal radiators are critical to maintain the lasers and spacecraft at a reasonable temperature. Our baseline is to keep the amplifiers near 300 K. The efficiency of the radiator can be determined by:

$$F = \dot{Q}/A = \epsilon\sigma T^4 \quad (1)$$

where ϵ is the emittance of the surface, σ is the Stefan-Boltzmann constant, T is the temperature, \dot{Q} is the heat rejected, A is the area, and F is the flux. The baseline radiator will be coated in AZ-93 white paint, which has a high emittance of 0.91 ± 0.02 (or conservatively, 0.89) and a low alpha, as it only absorbs 14-16% of incident sunlight on the spacecraft. The goal is to maintain a temperature of 300 K, as both the laser and onboard control electronics are operational at this temperature. At this temperature, the radiator can reject an idealized outward flux of 408 W/m². When taking into account the incident radiation, using a solar constant of 1,362 W/m² and a maximum 16% absorptance, the net flux of energy across the surface of the radiator is approximately 190 W/m². The baseline is to prevent direct solar illumination of the radiator.

The area of the radiator must be determined by thermal analysis, and is dependent on the desired operating temperature, heating from the environment, interactions with other surfaces of the spacecraft (*e.g.*, solar arrays), and the highest estimate (worst case) satellite waste heat. The waste heat in this case is dependent on the efficiency of the laser amplifiers—35%, 50%, or 70%, as mentioned. The worst-case estimate (35% efficiency) requires 65 kW to be rejected as waste heat for a 100 kW electrical input assuming virtually all the power goes to the laser (which is approximately correct during laser firing).

$$\dot{Q}_{\text{rejected}} = AF_{\text{net}} \quad (2)$$

where F_{net} is the net outward flux and \dot{Q} is the heat rejected. Given these parameters, the maximum required area of the radiator is ~341 m² for a 35% efficient laser amplifier. For a 50% efficient laser, a radiator area of ~262.1 m² is required; for a 70% efficient laser, a radiator area of ~157.2 m² is required.

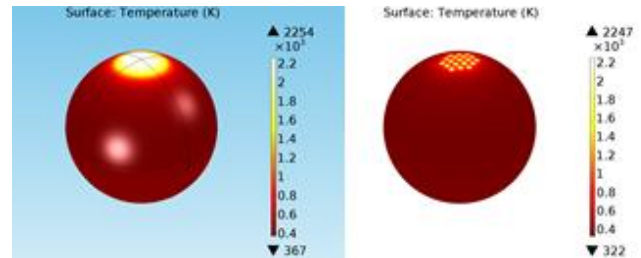


Figure 4. **Left:** Simulation showing one spot from the baseline phased array on the target at sufficient temperature to cause ablation. **Right:** Multi-beam simulation depicting 19 beams on the target from an optional choice of a close packed laser array instead of a phased array.

A passive cooling z-folded radiator consisting of two deployable panels will be used in order to provide a

sufficient surface area over which to emit the waste heat generated by the system. Each panel z-folds out into six segments, each of which further folds out into two additional segments, making 18 segments in total for each panel. The panels will rotate about their axes to maximize efficiency by remaining perpendicular to the sun and by radiating out of both sides. Each segment will be 2.2 m by 2.2 m, granting a total area of 348 m² out of which to dissipate heat. This will provide sufficient surface area to reject the maximum projected waste heat. If by the time of production, significant increases in laser efficiency have indeed been reached, the size of the panels can be altered so as to reduce the excess mass if less heat needs to be dissipated. Sun shades may also be employed to limit solar absorption and thus allow for greater efficiency. The current mass to power ratio for radiators is about 25 kg/kW for the ARM system as a baseline example. Radiators are currently the largest mass driver for large systems. This is an area that needs additional work, though even with the existing radiator designs, MW-class systems are feasible with current (or soon to exist) launchers. More laser amplifiers are easily added to allow for scaling to larger power levels. A 1 m to 4.5 m diameter is feasible; no additional deflection comes from the larger optic, just additional range from the target.

Launch Systems

The launch systems in consideration are Atlas V 551, Space Launch System (SLS) Block 1, Falcon Heavy, or Delta IV Heavy. These are likewise the launch systems in consideration for ARM, which calls for a payload of comparable parameters [13]. The DE-STARLITE spacecraft will fit within the payload fairing of any of the proposed launch systems (Fig. 5). As is evident from data in Table 1, the SLS Block 1 has the highest capabilities though a future design for the SLS Block 2 is projected to lift 130,000 kg to LEO. The Falcon Heavy has the smallest cost per unit mass, and has capabilities between that of the Atlas V and SLS Block 1. While the Atlas V 551 and Delta IV Heavy have previously undergone successful missions, the SLS Block 1 and Falcon Heavy are projected to be flight-proven within the timescale of the DE-STARLITE mission.

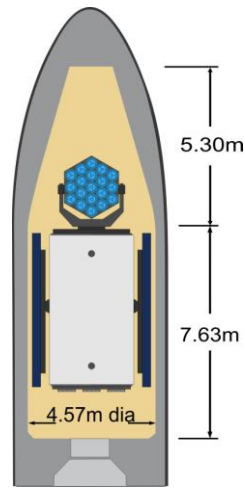


Figure 5. Stowed view of DE-STARLITE.

Table 1. Parameters of various launch vehicles in consideration for DE-STARLITE.

Parameter	Atlas V 551	SLS Block 1	Falcon Heavy	Delta IV Heavy
Payload Mass to LEO (kg)	18,500	70,000	53,000	28,790
Cost per unit mass to LEO	\$13 k/kg	\$19 k/kg	\$1.9 k/kg	\$13 k/kg
Fairing Diameter (m)	5.4	8.4	5.2	5
Status	Flight proven	Expected 2017	Expected 2015	Flight proven

As with ARM, it is possible to compensate for the lower capabilities of the Atlas V by using the SEP system to spiral out of Earth’s orbit and escape from Earth using Lunar Gravity Assist (LGA); however, this process of spiraling out and using LGA will take an additional 1 to 1.5 years of flight. All of these factors must be taken into consideration to choose the most effective launch system for the DE-STARLITE mission.

Launch mass

It is assumed in our analysis that a DE-STARLITE mission will use conventional launchers to get to LEO and then use ion engines, of a similar type to what is on the ARM mission to get from LEO to the target and then use the laser to do the actual target deflection. The launch mass is computed using many of the same assumptions used for the ARM mission but with the addition of the laser array and larger PV and radiators. The launch mass required vs. electrical power produced by the solar PV is shown. The laser power will be about 50% of the electrical power with assumed 50% conversion efficiency. This is a slightly optimistic assumption but well within the near term roadmap for the baselined Yb laser amplifiers.

Pointing and Control

The laser pointing control system uses a servo feedback based upon a SWIR camera observation of the spot intensity to control the phasing of each laser sub element to maximize the spot intensity. Phase control can also be used to move

the spot as needed on the target. Gross pointing is performed by the gimbal and bulk spacecraft motion. Spacecraft attitude control uses small ion engines.

4. ORBITAL DEFLECTION CAPABILITIES

This section describes how magnitude and duration of applied thrust influence miss distance. When an asteroid is exposed to the DE-STARLITE laser, the temperature (K) and flux (W/m^2) on the target asteroid must approach sufficiently high levels in order for significant ablation to occur, targeting a temperature on the order of 3,000 K and a flux of $>10^7 \text{ W}/\text{m}^2$. This causes direct evaporation of the asteroid at the spot of contact. Evaporation at the spot produces a vaporization plume thrust (N) that can be used to change the asteroid's orbit and effectively deflect asteroids from colliding with Earth. A miss distance of at least two Earth radii (12,742 km) is required to eliminate the threat of collision. The orbital deflection depends on the duration, magnitude, and direction of the applied thrust.

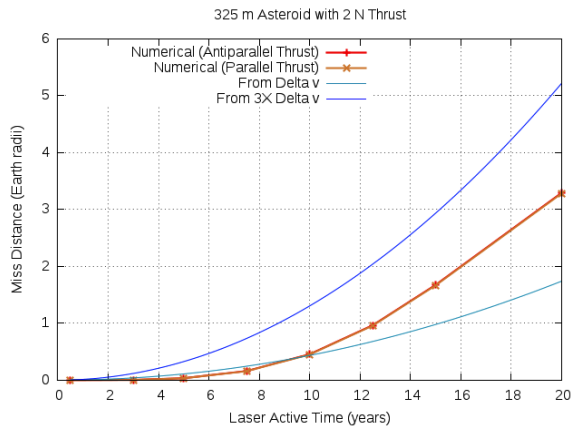


Figure 6. Miss distance vs. laser on time for orbital simulation with Δv and $3\Delta v$ approximations; nominal 2 N thrust ($\sim 30 \text{ kW}$ laser, a modest case for a DE-STARLITE mission). More thrust is available with larger arrays.

Full Numerical Orbital Simulations vs Analytic Approximations - A three-body simulation (accounting for the gravitational effects of the Earth, the sun, and the target asteroid) was performed in order to analyze how the applied thrust and the laser-active time impact the miss distance. In order to determine the orbital deflection, Δx , of an asteroid that is being acted on over a period of time, t , an approximation that is commonly used in orbital mechanics was used as a comparison. The detailed numerical simulation is compared to the approximation of multiplying by 3 the naive distance achieved by accelerating and coasting a system that is not a bound gravitational system. Hence the orbital deflection is compared to:

$$\Delta x_{\text{approx}} = 3(0.5 a \cdot t_{\text{active}}^2 + a \cdot t_{\text{active}} \cdot t_{\text{coast}}) \quad (3)$$

where a is the acceleration caused by the plume thrust, t_{active} is the time the laser is active, and t_{coast} is the coast time (typically zero). The reason this is done is because this approximation is often used for preliminary mission design. Fig. 6 compares the $1\Delta v$ and $3\Delta v$ approximations. A sample of the results for the 325 m asteroid case is shown

for the full numerical simulation of the orbital deflection along with the nominal Δv and $3\Delta v$ simplifications. It is evident that the $3\Delta v$ approximation is indeed only an approximation and in some cases fails badly.

The numerical simulations were performed in a rotating frame, where the thrust was pointed both along and against the velocity vector for comparison. Many dozens of orbital simulations were analyzed. Fig. 7 compares the laser-active time to the miss distance for a given thrust acting on targets of varying diameter. This focuses on the 325 m diameter asteroid case, as this is approximately the size of Apophis—a well-known possible threat. Computations have also been done for 20 to 1,000 m asteroids under many mission scenarios. The same code is used to analyze IBD, gravity tractor and impactor (impulse) cases to which DE-STARLITE are compared.

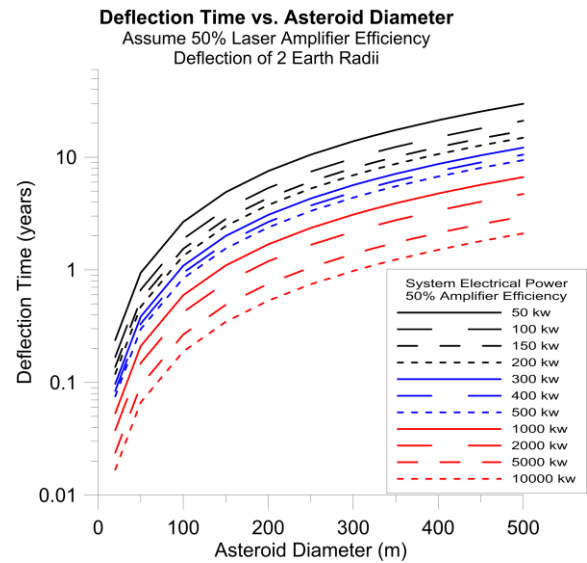


Figure 7. Estimated deflection time (laser on time) vs. target diameter and DE-STARLITE electrical power input from PV assuming a $3\Delta v$ approximation often over estimates the deflection (miss) distance. True mission planning requires detailed knowledge of the target orbit and the detailed interdiction scenario. Note that 200 m diameter asteroids can be deflected in ~ 1 year using a MW class laser; larger asteroids require more time. A MW laser DE-STARLITE mission appears to be launchable with an SLS Block 1.

5. IMPACTOR COMPARISON

Reference [1] discusses the case of IBD vs. laser ablation deflection. Here we discuss the case of using an impactor (ramming asteroid) vs. using a laser. As a common metric we use the launch mass as a common element for both cases—*i.e.*, for the same launch mass, what can each system do?

For a simplistic analysis the impactor delivers a large impulse or momentum transfer to deflect the target (integrated force - time in units of Ns). This momentum transfer imparts a change in the speed ΔV of the asteroid equals $\Delta p/M$ where M is the mass of the asteroid. Δp is the impulse delivered at a time τ before (if un-deflected) impact.

The term Δp equals mv where m is the spacecraft mass and v is the relative closing speed between the spacecraft and asteroid. The change of speed is thus

$$\Delta V = mv/M = v(m/M) \quad (4)$$

The deflection distance at the Earth is approximately

$$\Delta x = 3 \Delta V \cdot \tau = 3 \cdot v \cdot \tau (m/M) \quad (5)$$

where the factor of 3 is an approximation used from orbital dynamics but as we have shown in several of our papers it is not always a good approximation. We use it here for illustrative purposes and because it is often used in mission planetary defense planning exercises.

Note that the miss distance Δx is linearly proportional to the spacecraft mass (m), the closing speed (v) and time to impact τ and inversely proportional to the asteroid mass M . Note that the asteroid mass M is proportional to the cube of the asteroid diameter D . The momentum change (impulse delivered) is largely independent of the asteroid mass and only depends on the spacecraft mass (m) and the closing speed (v). For a homogeneous asteroid of density ρ then miss distance is:

$$\Delta x = 3 \Delta V \cdot \tau = 18 \cdot m \cdot v \cdot \tau / (\pi \rho D^3) \quad (6)$$

Since the asteroid is moving rapidly with typical speeds of 5-40 km/s we can simplify this to assume the spacecraft is simply in the way of the asteroid (inelastic billiard ball) and thus the speed of the spacecraft relation to the earth is of lesser importance. This of course depends on the specifics of the asteroid orbit (closing from the front vs. the back of the asteroid orbit). Essentially then it is the mass of the spacecraft that is critical to maximize. Once the spacecraft is launched to LEO it is assumed that ion engines will be used to allow a larger fraction of the launch mass to survive until impact to maximize the impulse. Since the miss distance is proportional to the inverse cube of the asteroid diameter, and the spacecraft mass is limited by the launcher capability, the only free parameter is the time to impact τ . Thus the miss distance is:

$$\Delta x = 3 \cdot \Delta V \cdot \tau = 3 \cdot \Delta p / M \cdot \tau = 18 \cdot m \cdot v \cdot \tau / (\pi \rho D^3) \quad (7)$$

In other words, the miss distance is proportional to:

$$\Delta x \sim m \cdot v \cdot \tau \cdot D^{-3} \quad (8)$$

For the case of directed energy the equivalent miss distance (using the same factor of 3 approximation for the effects of orbital mechanics) is:

$$\begin{aligned} \Delta x &= 3 \cdot 1/2 \cdot a \cdot \tau^2 = 3/2 (a \cdot \tau) \tau = 3/2 \Delta V \cdot \tau = 3/2 (F/M) \tau^2 \\ &= 3/2 F \cdot \tau^2 / M = 1/2 \cdot 3 \cdot \Delta p / M \cdot \tau = 9 \alpha P \cdot \tau^2 / (\pi \rho D^3) \end{aligned} \quad (9)$$

where:

- a = acceleration imparted due to the laser plume thrust
- F = laser plume thrust = αP
- P = laser power
- α = laser plume thrust coupling coefficient
- M = asteroid mass = $\pi \rho D^3 / 6$

We assume the laser thrust is constant and the asteroid mass changes very little due to the mass loss from ablation and that the laser plume thrust is proportional to the laser power. See our other papers on the detailed modeling for this. For simplicity we assume $\alpha \sim 80 \mu\text{N/W}$ optical in central spot. Note that for the case of directed energy or any constant

force (such as ion engines, gravity tractors, *etc.*) the miss distance:

$$\Delta x_{\text{laser}} = 1/2 \cdot 3 \cdot \Delta p / M \cdot \tau \quad (10)$$

while for the impulse delivery (effectively instantaneously at a time τ before impact) for the same overall delta momentum delivered to the asteroid is:

$$\Delta x_{\text{impactor}} = 3 \cdot \Delta p / M \cdot \tau, \text{ or: } \Delta x_{\text{laser}} = 1/2 \Delta x_{\text{impactor}} \quad (11)$$

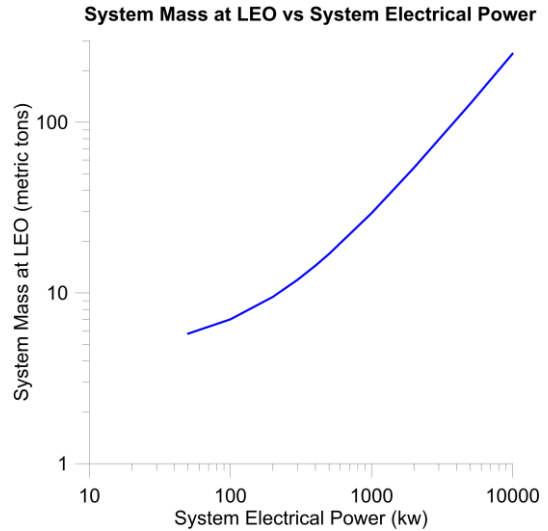


Figure 8. Mission mass at LEO vs. electrical power available from PV assuming nominal 50% laser amplifier efficiency, current ATM MegaFlex capability and $I_{sp} = 6,000$ s ion engines and radiator panels of 25 kg/kW radiated. SLS Block 1 launch of 70 metric tons to LEO corresponds roughly to 2-3 MW electrical or roughly 1 MW laser power.

Again this is for the simplistic assumption of the factor of 3 to approximate the orbital mechanics effects. The real situation is far more complex and depends on the specifics of the asteroid orbit and mission parameter as shown in Fig. 9 and Fig. 10. We assume an SLS Block 1 launch of 70,000 kg to LEO. For high I_{sp} ion engines of 3,000 s (Hall effect thrusters baselined for ARM) or 6,000 s (gridded ion) a decent fraction of the LEO mass will make it to the asteroid. For a comparable launch mass as would be needed for the Fig. 9 impactor case, if this same mass were used for the directed energy case, the laser exposure required would be about 1-2 years.

The details of the particular orbits are important but we can draw some basic conclusions. Assuming 60,000 kg makes it out to the asteroid and with a closing speed of 10 km/s, the impactor impulse is 6×10^8 Ns. Fig. 8 shows that for this same 70,000 kg SLS Block 1 to LEO, we could launch a 1 MW optical power laser delivering ~ 60 N of thrust on the asteroid for an assumed laser coupling coefficient $\alpha \sim 80 \mu\text{N/W}$ optical with an assumed (somewhat optimistic) high efficiency beam formation in the central spot of 0.7. To get the same deflection in the same time to impact as the impactor, we need the laser system to deliver twice the momentum as the impactor. Hence, we need 1.2×10^9 N s. At 60 N of laser plume thrust this would require a time $\tau = 1.2 \times 10^9$ N s / 60 N = 2×10^7 s or about 7

months. In this case, the exposure time needed is about 7 months. This time is independent of the launch mass as both the impactor momentum delivered and the laser momentum are proportional to launch mass for reasonably large launch masses. Other differences for real systems are typical impactor missions need more than one to make sure the impulse was delivered properly. For any real threat, multiple backups would be prudent.

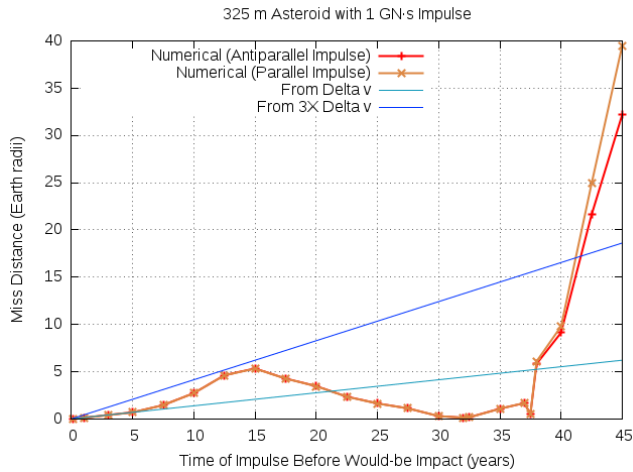


Figure 9. Miss distance vs. impulse delivery time before impact for 1 GN s impulse (325 m asteroid). This is somewhat larger than an SLS Block 1. A miss distance of 2 Earth radii (typ. min acceptable) would require interdiction about 10 years before impact. The seemingly unusual behavior from the full simulation is due to resonance effects from the multiple orbits. It is clear the $3\Delta v$ approximation is not always accurate, and can be very misleading in some cases.

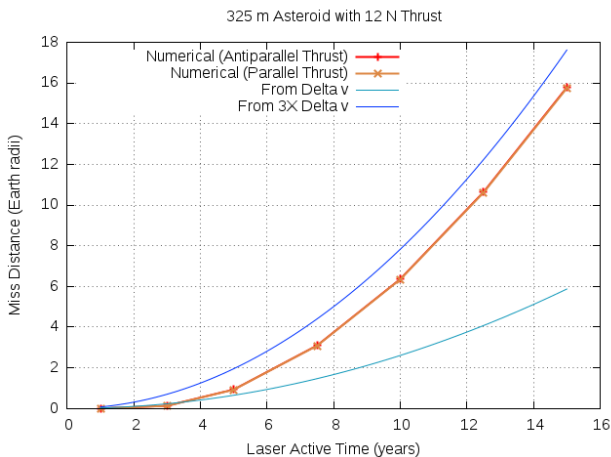


Figure 10. Miss distance vs. laser exposure time for 12 N thrust on a 325 m diameter asteroid. Parallel and anti-parallel cases are coincident in the plot. An SLS Block 1 could deliver $\sim 5x$ this thrust. A 2 Earth radii miss requires ~ 6 years of exposure.

6. ION BEAM DEFLECTION COMPARISON

Ion beam deflection (IBD) is an alternative approach to achieve asteroid orbital deflection in which an ion beam is used to push against the asteroid. In using this approach, the spacecraft must provide twice as much thrust

as would otherwise be necessary to deflect the asteroid a desired distance. Half of the thrust is lost in station keeping in order to keep the spacecraft stable, as the spacecraft must push towards or away from the asteroid with an equal amount of thrust. This comparison is discussed in detail in [13], [19]. The basic issue is that in order for an IBD mission to be effective against a large asteroid it must carry a large amount of ion propellant (currently Xe) and the required deflection propellant scales as the mass of the asteroid or as D^3 where D is the asteroid diameter. An example comparing the launch mass of an IBD to laser deflection mission is shown in Figure 11. This clearly shows the advantage of the laser deflection mission. For an equivalent warning time, the IBD case with an I_{sp} of 3000 s requires ~ 125 kW electrical power, and the IBD case with an I_{sp} of 6000 s requires ~ 250 kW electrical power. The same parameters (8.5 year build and travel time, 50% efficient laser amplifiers, 2 g/cc and 2 Earth radii miss distance) are assumed. Note that the 8.5 year build and travel time is assumed for a spacecraft using ion engines with an I_{sp} of 3000 s; the travel time (typ. $\sim 1-2$ year) may be decreased with ion engines of greater specific impulse and efficiency. Build time can be reduced to essentially zero with pre-deployed missions. If we pre-deploy the deflection assets the warning time is reduced by about 5 years (typ minimum “start mission design” to launch).

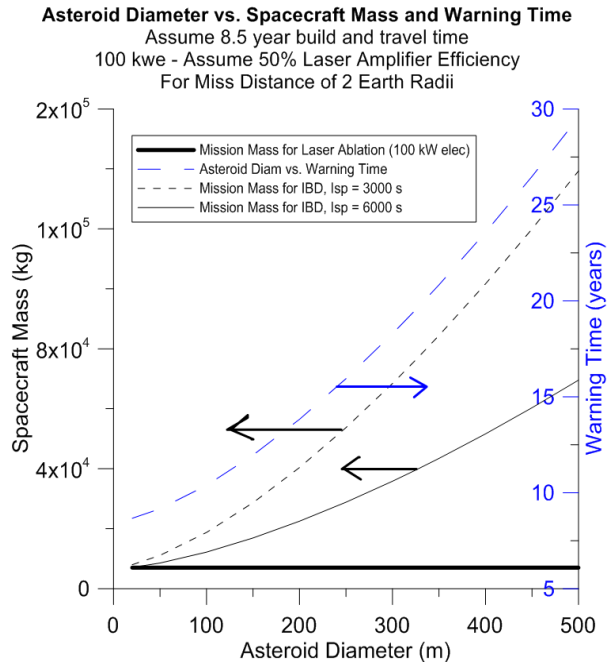


Figure 11. Asteroid diameter vs. spacecraft mass at LEO (left axis) for the IBD case (magnetically shielded Hall effect thrusters w/ I_{sp} of 3000 s, and gridded ion thrusters w/ I_{sp} of 6000 s) and for laser ablation, as well as asteroid diameter vs. the required warning time for a modest laser ablation system with 100 kW electrical power (right axis).

7. GRAVITY TRACTOR COMPARISON

Gravity tractors are potentially very attractive due to their simplicity. They simply act as a gravitational bound “tug” whose ultimate deflection is due to the propulsion of

the spacecraft. The gravitational attraction between the spacecraft and the asteroid is simply a “gravity tug rope” with a very low strength set by the gravitational attraction between the two objects. An enhanced gravity attractor is similar except additional mass is added to the spacecraft by removing it from the asteroid and transferring it to the spacecraft. The actual deflection is accomplished by using the propulsion of the spacecraft to “pull” the asteroid and thus deflect it. In this sense it is the same as the IBD approach EXCEPT there is not an ion beam pushing the asteroid but rather the ion beam is used for propulsion of the spacecraft (as it is for the IBD case) and this propulsion thus pulls on the asteroid (via the gravity coupling between the spacecraft and asteroid). The ion beam propulsion must be less than the gravitation attraction “gravity tug rope strength” or else the spacecraft will simply escape from the asteroid rather than pull it. The enhanced gravity tractor simply increases the “gravity tug rope strength” but does not affect the total ion beam engine delta P. The advantage of this over IBD is roughly a factor of two less ion engine fuel is needed since no ion beam fuel is use to push the asteroid (as well as an equal and opposite beam to station keep the IBD mode. In this way gravity tractors use ½ the fuel (typ Xe gas) for the (ion) propulsion as the IBD case BUT the disadvantage is that the “gravity tug rope strength) is very weak due to the weak gravitational attraction between the spacecraft and the asteroid. This effectively means the required deflection force (equal and opposite to the gravity attraction) is extremely low thus requiring an extremely long time to deflect the asteroid.

For example if we have an asteroid of density ρ , radius R and mass M with a spacecraft of mass m (whether an enhanced gravity tractor or not) then the force between the spacecraft and asteroid and hence the “gravity tug rope strength” is: $F = GMm/(R+s)^2$ where s is the distance the spacecraft is from the surface and $G \sim 6.67 \times 10^{-11}$ (MKS). Here $M = \rho 4/3\pi R^3$ or $F = G \rho 4/3\pi m R^3 / (R+s)^2$. The maximum this can be is when the spacecraft is at the surface or $s=0$. While not realistic this is the maximum force we can impart to the asteroid by pulling on it using the spacecraft engines as this is the max force before “breaking” the gravitational bond between the spacecraft and asteroid. For $s=0$ we get $F = G \rho 4/3\pi m R$. Note $F \sim Rm$ and hence larger forces are possible (stronger “gravity tug rope”) for larger asteroids. For typ $\rho \sim 2000$ kg/m³ we get F (μ N) $\sim 0.56 R(m)$ m(kg). For $R=100$ m and a spacecraft mass $m=10^4$ kg we get $F = 0.56$ N. We can increase the coupling by increasing the spacecraft mass (expensive in terms of launch systems and ion propellant required to get the spacecraft to the asteroid,) or using the asteroid itself to transfer mass to the spacecraft once it arrives. This is the “enhanced gravity tractor”. Assuming this can be done we can then increase the spacecraft mass m at the expense of the asteroid mass M. We will ignore the complexity and energetic of doing this and thus analyze the most optimistic case for the enhanced gravity tractor. The net effect of transferring asteroid mass to the spacecraft is to increase m and thus increase the gravitational attraction and thus increase the “gravity tug rope strength” or F. Note that the effective asteroid radius R

does not change dramatically even if we transfer ½ the asteroid mass (an extreme case) since the asteroid mass scales as $M \sim R^3$ and hence $R \sim M^{1/3}$ thus even transferring ½ the mass (reducing M by a factor of 2) only reduces R by 1.26. However, this does allow the enhanced gravity tractor to increase the maximum tug force possible to much larger values allowing us to use larger spacecraft engines (currently ion engines). This is impressive but unfortunately it does nothing to deflect the asteroid. To deflect the asteroid we must use the spacecraft engines. Thus we are back to the equivalent IBD case BUT with ½ the ion fuel load. Referring to Fig 11 we see that the launch mass for the IBD (and thus similar for gravity tractor) cases increases rapidly with the asteroid diameter and the deflection time also increases compared to the same mass for DE-STARLITE. In the case of the normal (unenhanced) gravity tractor, the system is largely not tenable due to the very weak gravity coupling.

Gravitation Escape Speed and Disassembly - If we go to the bother of landing on the asteroid and transferring material for an enhanced gravity tractor then it might be better to propel the material off the surface at a speed that exceeds the gravitational escape speed. When synchronized properly this is the most energetically efficient solution to deflection. The escape speed for a body of uniform density ρ (kg/m³) with radius R (m) is:

$$v_{esc} = \sqrt{\frac{8\pi}{3} G \rho R} \text{ or about } v_{esc} \approx 2.36 \times 10^{-5} \sqrt{\rho R}$$

As an example for $R=100$ (200 m diameter) and $\rho=2000$ kg/m³ we get $v_{esc} \sim 0.11$ m/s. The total gravitational binding energy of the asteroid is:

$$G_{BE} = 3/5 \frac{GM^2}{R} \text{ or assuming constant density:}$$

$$G_{BE} = \frac{16}{15} G \pi^2 \rho^2 R^5 \sim 7.02 \times 10^{-10} \rho^2 R^5$$

For $R=100$ (200 m diameter) and $\rho=2000$ kg/m³ we get:

$$G_{BE} \sim 2.81 \times 10^7 \sim 28 MJ .$$

Note that 1 ton TNT has an energy of about 4GJ. Thus our 200m diameter asteroid has a gravitational binding energy of about 6.5 kg (14 lbs) of TNT! One way of dealing with a asteroid threat is to use an array (MIRV) of penetrating explosive and non explosive impactors to gravitational debind it. The primary issue will be whether molecular binding will be dominant or not (ie loose rock pile or solid) and the scale of disassembly (size distribution of debris). If hit early enough the broken asteroid material will miss or if hit late the atmosphere then becomes a “beam dump”. It is important to get the typical debris size to the meter scale in the latter case.

8. PRE-DEPLOYMENT OF DEFENSE ASSET

No terrestrial defense asset has a “build start” upon detection of the threat. This would clearly be ludicrous. No one would imagine starting to build a missile defense interceptor AFTER the warhead launch is detected. This would be considered ridiculous. It should be no different for

planetary defense. We propose pre-deployment of the defensive systems BEFORE the threat is recognized just as we do with terrestrial defense. This is the only practical solution to planetary defense systems. We are used to building terrestrial defense system with the hope that we will NEVER use them. The same attitude should apply to planetary defense – even more so. On one SLS Block 1 (70 metric tons to LEO) we can fit a number of DE-STARLITE systems (Fig 7 and 8) that could be deployed at LEO or boosted to GEO awaiting a threat. An often quoted argument is “why build a system before a threat is identified”? It is a waste of money. This same argument is never used for terrestrial threats. You do not hear “why build a missile defense system against rogue states”. Wait until we detect a launch then we will build it! Of course not. So why do we accept this argument for planetary defense. We should not. **Pre-deployment is the only logical solution.** A full defense system must clearly include significantly enhanced detection assets.

9. CONCLUSIONS

Directed energy for planetary defense is a very promising planetary defense system at a modest cost. As outlined above, DE-STAR and DE-STARLITE employ laser ablation technologies which use the asteroid as the "fuel" for its own deflection. In particular, DE-STARLITE is able to mitigate much larger targets than would be possible with other proposed technologies such as IBD, gravity tractors, and kinetic impactors. For instance, with the equivalent mass of an ARM Block 1 arrangement (14 tons to LEO), designed to capture a 5-10 m diameter asteroid, DE-STARLITE can mitigate an asteroid larger than Apophis (325 m diameter), even without keyhole effects. A full SLS block 1 (70 tons to LEO) is capable of deploying a number of DE-STARLITE systems simultaneously. Much smaller DE-STARLITE systems could be used for testing on targets that are likely to pass through keyholes. The same technology proposed for DE-STARLITE has significant long-range implications for space missions, as outlined in other DE-STAR papers. Among other benefits, the DE-STARLITE system utilizes rapidly developing technologies to perform a task previously thought to be mere science fiction and can easily be increased or decreased in scope given its scalable and modular nature. DE-STARLITE is capable of launching on an Atlas V 551, Falcon Heavy, SLS, Ariane V or Delta IV Heavy, among others. Many of the items needed for the DE-STARLITE system currently have high technology readiness level (TRL); however, one critical issue currently being worked on is the radiation hardening of the lasers, though it appears achievable to raise this to a TRL 6 within 3-5 years. Laser lifetime also poses an issue, though this is likewise being worked on; a path forward for continuous operation looks quite feasible, with or without redundancy options for the lasers. Given that the laser amplifier mass is small and the system is designed to take multiple fibers in each configuration, redundant amplifiers can be easily implemented if needed. DE-STARLITE is a critical step

towards achieving the long-term goal of implementing a standoff system capable of full planetary defense and many other tasks including spacecraft propulsion. DE-STARLITE represents a practicable technology that can be implemented within a much shorter time frame at a much lower cost. DE-STARLITE will help to establish the viability of many of the critical technologies for future use in larger systems. We strongly propose pre-deployment of the systems as a practical measure for planetary defense rather than waiting for the threat before building. We would never accept such an attitude for terrestrial missile defense, nor should we for planetary defense.

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