Effects of Matter

Matter is polarized by both electric and magnetic fields. The dipoles (both electric and magnetic) create their own additional E and B fields which add (or subtract) to externally applied field to form a total field that is modified from the applied field.

Maxwell's Equations of Electromagnetism

No different in matter than in vacuum BUT E and B are the TOTAL E,B

 ρ is the total charge density and \vec{J} is the total current density - these include internal effects of matter

1)
$$\nabla \bullet \vec{E} = \rho / \varepsilon_0 \ (\rho = \text{charge/volume or Coulomb/m}^3) \quad \oint_{\Sigma} \vec{E} \bullet d\vec{A} = Q_{enc} / \varepsilon_0$$

2) $\nabla \bullet \vec{B} = 0 \qquad \qquad \oint_{\Sigma} \vec{B} \bullet d\vec{A} = 0$

(no magnetic monopoles found yet - if they exist then $\nabla \bullet \vec{B} \propto \rho_m$ (Mag monopole density)

3)
$$\nabla x \vec{E} = -\frac{\partial B}{\partial t}$$

4) $\nabla x \vec{B} = \varepsilon_0 \mu_0 \frac{\partial \vec{E}}{\partial t} + \mu_0 \vec{J}$
(\vec{J} = current density vector (amp/m²) $I = \int_{\Sigma} \vec{J} \cdot d\vec{A}$)
Define Displacement Current as:

$$\vec{J}_{D} \equiv \varepsilon_{0} \frac{\partial \vec{E}}{\partial t}$$

$$4) \nabla x \vec{B} = \varepsilon_{0} \mu_{0} \frac{\partial \vec{E}}{\partial t} + \mu_{0} \vec{J} = \mu_{0} (\vec{J}_{D} + \vec{J}) \qquad \qquad \oint_{C} \vec{B} \bullet d\vec{l} = \mu_{0} I_{D} + \mu_{0} I_{enc}$$

Note that \vec{J} is real current density while \vec{J}_D is a ficticious current density Note that Stokes Theorem for ANY well behaved vector field \vec{F} : $\oint \vec{F} \bullet d\vec{l} = \int \nabla x \vec{F} \bullet d\vec{A}$

EM Wave Equation in Materials

 \rightarrow Note the matters terms $\frac{\partial J}{\partial t} + \nabla(\rho_e / \varepsilon_0)$ and $\nabla x \vec{J}$

E field

 $\nabla x(\nabla x \vec{E}) = \nabla (\nabla \bullet \vec{E}) - \nabla^2 (\vec{E}) = \nabla (\rho_e / \varepsilon_0) - \nabla^2 (\vec{E})$ $= \nabla x(-\frac{\partial \vec{B}}{\partial t}) = -\frac{\partial (\nabla x \vec{B})}{\partial t} = -\frac{\partial (\varepsilon_0 \mu_0 \frac{\partial \vec{E}}{\partial t} + \mu_0 \vec{J})}{\partial t} = -\varepsilon_0 \mu_0 \frac{\partial^2 \vec{E}}{\partial t^2} - \mu_0 \frac{\partial \vec{J}}{\partial t}$ $\rightarrow \nabla^2 (\vec{E}) - \varepsilon_0 \mu_0 \frac{\partial^2 \vec{E}}{\partial t^2} = \mu_0 \frac{\partial \vec{J}}{\partial t} + \nabla (\rho_e / \varepsilon_0)$

B field

 $\nabla x(\nabla x \vec{B}) = \nabla (\nabla \bullet \vec{B}) - \nabla^2 (\vec{B}) = \nabla (0) - \nabla^2 (\vec{B})$ $= \nabla x(\varepsilon_0 \mu_0 \frac{\partial \vec{E}}{\partial t} + \mu_0 \vec{J}) = \varepsilon_0 \mu_0 \frac{\partial (\nabla x \vec{E})}{\partial t} + \mu_0 \nabla x \vec{J} = -\varepsilon_0 \mu_0 \frac{\partial^2 \vec{B}}{\partial t^2} + \mu_0 \nabla x \vec{J}$ $\rightarrow \nabla^2 (\vec{B}) - \varepsilon_0 \mu_0 \frac{\partial^2 \vec{B}}{\partial t^2} = -\mu_0 \nabla x \vec{J}$

Polarized matter

No free charges in this case We will allow free flow later

- E,B dipoles are formed by applied EM wave
- EM wave has time varying (oscillating) E,B
- Oscillating E,B \rightarrow induced oscillating dipoles in matter
- Oscillating dipoles \rightarrow accelerated charges
 - Can think if time varying dipole moments
- Oscillating charges \rightarrow radiation
- Radiation is EM wave \rightarrow adds/ subtracts
 - From applied EM wave
 - \rightarrow new EM wave is sum of the applied and induced wave
- It is this effect (polarization) that is the basis of most optical properties you will deal with
- While BOTH E and B dipoles are allowed most materials are dominated by the effect of electric dipole moments
 - Ex: "non magnetic doped" glass, plastic, water, air, ...

Electric Dipoles in Matter

Consider neutral atoms and molecules that can be polarized

- $\overline{E_a}$ = atomic/molecular E field induced by local (nearby) E field
- $\overrightarrow{E_L}$ = local electric field due to all other polarized matter + external field

For class A dielectrics (most of the materials you deal with): $\overrightarrow{E_a} \propto \overrightarrow{E_L}$ BUT opposes it

 $\overrightarrow{E_T} = \overrightarrow{E_L} + \overrightarrow{E_a} \text{ total electric field DUE to BOTH the polarized local matter field + rest of matter + external field}$ $\left\langle \overrightarrow{E_T} \right\rangle = \left\langle \overrightarrow{E_L} \right\rangle + \left\langle \overrightarrow{E_a} \right\rangle \text{ what we measure in matter}$ $\text{Note: } \left| \left\langle \overrightarrow{E_L} \right\rangle \right| > \left| \left\langle \overrightarrow{E_T} \right\rangle \right| \text{ since } \overrightarrow{E_a} \text{ opposes } \overrightarrow{E_L}$

Consider a uniform spherical charge distribution with density ρ

$$\oint_{\Sigma} \vec{E} \bullet d\vec{A} = 4\pi r^2 E(r) = Q_{enc} / \varepsilon_0 = \frac{4\pi}{3} r^3 \rho \to E(r) = \frac{r\rho}{3\varepsilon_0}$$

 $\vec{E}(r)$ is radial from center of charge distribution

Atomic and Molecular Dipole

Two charges +-q set a distance vector apart of \vec{l} Distance vector from +q to obs point is \vec{r} ' Distance vector from -q to obs point is \vec{r} Potenital $\phi(\mathbf{r}) = \operatorname{kq}(\frac{1}{r} - \frac{1}{r}) \ k = \frac{1}{4\pi\varepsilon_0}$ $\vec{r} = \vec{r}' + \vec{l} \quad \vec{r}' = \vec{r} - \vec{l}$ $r' = |\vec{r}'| = [(\vec{r} - \vec{l}) \cdot (\vec{r} - \vec{l})]^{1/2} = [r^2 + l^2 - 2\vec{r} \cdot \vec{l}]^{1/2} = r[1 + l^2 / r^2 - 2\vec{r} \cdot \vec{l} / r^2]^{1/2}$ Assume distances r,r' are large compared to atomic distance l $l^2 / r^2 \sim 0$ $\Rightarrow r' \approx r[1 - 2\vec{r} \cdot \vec{l} / r^2]^{1/2} \approx r(1 - \vec{r} \cdot \vec{l} / r^2)$

$$\rightarrow \frac{1}{r'} \sim \frac{1}{r} (1 + \vec{r} \cdot \vec{l} / r^2) \rightarrow \frac{1}{r'} - \frac{1}{r} \sim \vec{r} \cdot \vec{l} / r^3)$$
$$\rightarrow \phi(\mathbf{r}) = kq(\frac{1}{r'} - \frac{1}{r}) \sim k \frac{q\vec{r} \cdot \vec{l}}{r^3}$$

Define $\vec{p} = q\vec{l}$ = electric dipole moment

$$\phi(\mathbf{r}) \sim k \frac{\vec{\mathbf{p}} \bullet \vec{r}}{r^3} = k \vec{\mathbf{p}} \bullet \frac{\vec{r}}{r^3} = -k \vec{\mathbf{p}} \bullet \nabla(1/r)$$

Recall $\nabla(1/r) = -\frac{\vec{r}}{r^3}$ (spherical gradient

Electric Field from Atomic Dipole

 $\overrightarrow{E_a}$ = atomic/molecular E field induced by local (nearby) E field

$$\overrightarrow{E_a} = -\nabla(\phi) = \nabla(k\overrightarrow{p} \bullet \nabla(1/r)) \quad (r >> l)$$

Average over volume V

$$\left\langle \overrightarrow{E_a} \right\rangle = \frac{1}{V} \int \overrightarrow{E_a} dV = \frac{k}{V} \int \nabla(\overrightarrow{p} \bullet \nabla(1/r)) dV = \frac{k}{V} \oint_{\Sigma} \overrightarrow{p} \bullet \nabla(1/r) d\overrightarrow{A}$$

 $\operatorname{Recall} \int \nabla(f) dV = \oint_{\Sigma} f d\vec{A}$

Let the volume V be a sphere of radius b centered on a microscopic dipole

 $\nabla(1/r) = -\frac{r}{r^3} = -\frac{r}{h^3}$ (at surface of sphere) = $-\frac{n}{h^2}$ (where \hat{n} is the radial outward normal vector) $\left\langle \vec{E_a} \right\rangle = \frac{k}{V} \oint \vec{p} \bullet \nabla (1/r) d\vec{A} = \frac{-k}{b^2 V} \oint \vec{p} \bullet \hat{n} d\vec{A} = \frac{-k}{b^2 V} \frac{1}{3} 4\pi b^2 \vec{p} = -\frac{p}{3\epsilon V}$ Can show $\oint \vec{p} \cdot \hat{n} d\vec{A} = \oint \vec{p} \cdot \hat{n} \hat{n} dA = \frac{1}{3} 4\pi b^2 \vec{p}$ $\hat{n} = [\sin(\theta)\cos(\phi), \sin(\theta)\sin(\phi), \cos(\theta)], \text{ let } \vec{p} = (0, 0, p) (z \text{ axis } p)$ $dA = b^2 \sin(\theta) d\theta d\phi$ $\vec{N} = \oint_{n} (\vec{p} \cdot \hat{n}) \hat{n} dA = \oint_{n} p \cos(\theta) \hat{n} dA = \oint_{n} p \cos(\theta) \hat{n} b^{2} \sin(\theta) d\theta d\phi = p b^{2} \oint_{n} \hat{n} \cos(\theta) \sin(\theta) d\theta d\phi$ $N_{x} = N_{y} = 0$ $N_z = pb^2 \oint \cos^2(\theta) \sin(\theta) d\theta d\phi = 2\pi pb^2 \int_0^{\pi} \cos^2(\theta) \sin(\theta) d\theta = -2\pi pb^2 \frac{\cos^3(\theta)}{3} l_0^{\pi} = \frac{4\pi}{3} pb^2$

(note error in notes - extra factor of b on Physics 141 - Notes 2.pdf)